

Chromospheric emission from nanoflare heating in RADYN simulations

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ABSTRACT

Context. Heating signatures from small-scale magnetic reconnection events in the solar atmosphere have proven to be difficult to detect through observations. Numerical models that reproduce flaring conditions are essential in understanding how nanoflares may act as a heating mechanism of the corona.

Aims. We study the effects of non-thermal electrons in synthetic spectra from 1D hydrodynamic RADYN simulations of nanoflare heated loops to investigate the diagnostic potential of chromospheric emission from small-scale events.

Methods. The Mg II h and k, Ca II H and K, Ca II 854.2 nm, and H α and H β chromospheric lines were synthesised from various RADYN models of coronal loops subject to electron beams of nanoflare energies. The contribution function to the line intensity was computed to better understand how the atmospheric response to the non-thermal electrons affects the formation of spectral lines and the detailed shape of their spectral profiles.

Results. The spectral line signatures arising from the electron beams highly depend on the density of the loop and the lower cutoff energy of the electrons. Low-energy (5 keV) electrons deposit their energy in the corona and transition region, producing strong plasma flows that cause both redshifts and blueshifts of the chromospheric spectra. Higher-energy (10 and 15 keV) electrons deposit their energy in the lower transition region and chromosphere, resulting in increased emission from local heating. Our results indicate that effects from small-scale events can be observed with ground-based telescopes, expanding the list of possible diagnostics for the presence and properties of nanoflares.

Key words. Sun: chromosphere – Sun: flares – Methods: numerical – Radiative transfer – Line: profiles

1. Introduction

Nanoflares are small-scale events associated with magnetic reconnection in the solar atmosphere. The nanoflare heating mechanism is one of the prime candidates in understanding the high temperature of the corona (Parker 1988). Flare energy is believed to be transported by electrons from the thermal background, where the electrons are accelerated to non-thermal energies as the magnetic field lines reconnect. The accelerated electrons travel along the magnetic field through the transition region (TR) and chromosphere, where the energy is lost via Coulomb collisions with the ambient plasma (Brown 1971; Emslie 1978; Holman et al. 2011). This gives rise to observable signatures in spectral lines that are formed in the sites where energy is deposited. However, spectral line signatures are also affected by the response to the dissipated electron energy elsewhere in the atmosphere, such as strong plasma flows.

Signatures of non-thermal electrons are found in spectra from hard X-ray observations of active region flares. However, X-ray observations of small-scale events with energies in the range 10^{24} – 10^{25} erg remain rare because signatures from non-thermal electrons are typically below the detection threshold (although, see e.g. Wright et al. 2017; Glesener et al. 2020; Cooper et al. 2021). Indications of energetic events are obtained in regions that are responsive to heating. In the corona, such indica-

tions are more difficult to observe because its high conductivity smears the heating signatures out. The TR, on the other hand, is highly responsive to heating because of rapid changes in density, temperature, and volume during heating events. Non-thermal electrons from coronal heating events give rise to changes in density and temperature through collisions with the dense TR and chromospheric plasma. Through advanced numerical simulations, Testa et al. (2014) have demonstrated that blueshifts in the Si IV 140.2 nm line observed with the Interface Region Imaging Spectrograph (IRIS; De Pontieu et al. 2014) in small heating events at the footpoints of transient hot loops are consistent with heating by non-thermal electrons, and they cannot be reproduced by models that assume heating only by thermal conduction. Polito et al. (2018) carried out an extensive numerical investigation to better understand and interpret TR observations by exploring a wide region of parameters. Testa et al. (2020) have further shown that observations of high variability (≤ 60 s) at the footpoints of active region (AR) high-temperature coronal loops (~ 8 – 10 MK) in combination with numerical simulations provide powerful diagnostics of coronal heating and energy transport.

Polito et al. (2018) analysed in detail the TR Si IV lines and upper chromospheric Mg II h and k lines from RADYN 1D simulations of nanoflare heated loops, and Testa et al. (2020) also investigated C II and Mg II triplet emission from the same set of simulations. In this work, we extend their analyses by probing

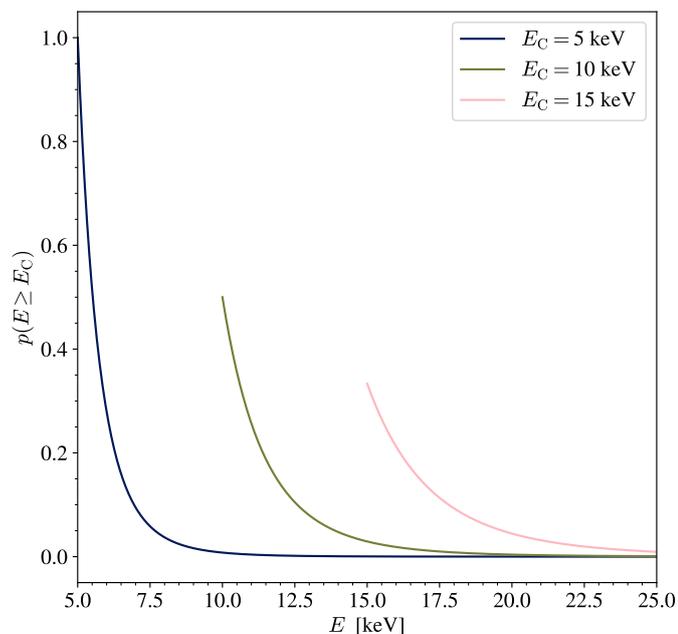


Fig. 1. Sketch of power-law distributions with different cutoff energies E_C , on a linear scaling. The figure serves as a visual aid to distinguish the difference between the electron distributions.

the impact of electron beam heating on spectral lines that are formed deeper in the atmosphere and are readily accessible by ground-based telescopes: Ca II H and K, Ca II 854.2 nm, and H α and H β . We focus our efforts on the electron beam heating models with the shorter loop length (15 Mm) and initial loop-top temperatures of 1 and 3 MK, representing respectively an initially almost empty loop and a typical AR core loop. The selection of simulations is motivated by the study of nanoflare electron beams in 3D MHD Bifrost models (Bakke et al. 2018; Frogner et al. 2020), which only covers cooler and shorter loops. With RADYN we can more easily study the details of such conditions, while with Bifrost this is more difficult due to the complex configuration of the simulation. We reproduce the upper chromospheric Mg II h results of Polito et al. (2018) with the aim of making a direct comparison with the lower chromospheric observables. In the analysis we explore the impact of varying the low-energy cutoff value by investigating the atmospheric response, the Doppler shift of the chromospheric lines, and the formation of line intensity.

2. Method

2.1. The RADYN numerical code

The numerical simulations were performed using the RADYN numerical code (Carlsson & Stein 1992, 1995, 1997; Allred et al. 2015). RADYN solves the equation of charge and population conservation coupled to the non-linear equations of hydrodynamics in a 1D plane-parallel atmosphere on an adaptive grid (Dorfi & Drury 1987). The non-local thermodynamic equilibrium (non-LTE) radiation transport is solved for H, He, and Ca II, while continua from other atomic species are treated in local thermodynamic equilibrium (LTE) as background metal opacities using the Uppsala opacity package (Gustafsson 1973). Optically thin losses are calculated using the CHIANTI 7.1 atomic database (Dere et al. 1997).

The RADYN version described in Allred et al. (2015) allows to model the effect of non-thermal electrons in flaring loops. The particle distribution function is calculated by the Fokker-Planck equations, and the electron beam spectrum follows a power-law. With RADYN flare models it is possible to study the atmospheric response to energy deposited in the loop. RADYN models half of the loop only, between one footpoint and the loop apex. It is implicitly assumed that the full loop is symmetric around the apex. Further, the half-loop is modelled as a quarter circle, where the geometry is used to vary the value for the gravitational acceleration along the z -axis. The radiative transfer is computed by assuming a 1D atmosphere along the z -axis. The spectral diagnostics of interest are formed in the lowest part of the atmosphere where the loop geometry is near-vertical. This allows for comparisons with observations with a top-down line of sight.

We have investigated the heating models of coronal nanoflares carried out by Polito et al. (2018), but focused our efforts on the models with transport by non-thermal electrons. The atmospheric structure included a plage-like chromosphere following the work by Carlsson et al. (2015). The numerical investigation included several nanoflare models with different loop-top temperatures and parameters of the non-thermal electron power-law distribution.

2.2. Parameter survey

A total of six numerical simulations were analysed. The energy deposited in a single loop was set to $E = 6 \cdot 10^{24}$ erg in all models, following the work by Testa et al. (2014, 2020). The cross-sectional area of the loop was set to $A = 5 \cdot 10^{10}$ m², corresponding to a diameter of about 250 km. The half-loop lengths were set to 15 Mm with initial loop-top temperatures $T_{LT} = 1$ and 3 MK, giving two initial atmospheres that aimed to reproduce AR loops at different heating stages. With the assumption of an atmosphere in hydrostatic equilibrium, low-temperature loops have low density while hotter loops are also denser. The 1 MK loop represents an empty strand where heating has not yet occurred, and has an apex electron number density around $10^{8.7}$ cm⁻³. The 3 MK loop represents a previously heated strand with apex electron number density of approximately $10^{9.6}$ cm⁻³. In the following, we refer to the 1 MK and 3 MK loops as empty and pre-heated loops, respectively.

In RADYN flare models, the energy flux F , spectral index δ , and low-energy cutoff E_C that go into the power-law distribution are user-specified. The input value for the electron beam energy flux was calculated as

$$F = \frac{E}{At} = 1.2 \cdot 10^9 \text{ erg cm}^{-2} \text{ s}^{-1}, \quad (1)$$

where E and A are the total energy and cross-section of the loop given above. The loop was heated constantly for $t = 10$ s which is consistent with IRIS observations of TR moss showing that the lifetime of short-lived footpoint brightenings varies between 10 – 30 s (Testa et al. 2013, 2014, 2020).

A large spectral index ($\delta = 7$) was chosen for the energy distribution, supported by observational evidence of increasing spectral index with decreasing flare energy (Hannah et al. 2011), and compatible with values from rare hard X-ray observations of these small heating events (e.g. Glesener et al. 2020; Cooper et al. 2021). The low-energy cutoff values $E_C = 5, 10,$ and 15 keV were chosen to represent beams that are dominated by low-, intermediate-, and high-energy electrons. Figure 1 shows examples of power-law distributions for the different values of

E_C , where the total energy of the beam is fixed. Each distribution starts at its respective cutoff energy (which is the energy that most of the electrons have because of the large spectral index), before exponentially declining as the energy increases. The $E_C = 5$ keV distribution shows that there are few electrons with $E > 9$ keV, and that most of the electrons have energies between 5 and 6 keV. As the cutoff energy increases, the number of electrons with $E = E_C$ decreases. However, there are generally more electrons with higher energies in these distributions.

2.3. Spectral synthesis with RH

The chromospheric lines were synthesised using the RH1.5D radiative transfer code (Uitenbroek 2001; Pereira & Uitenbroek 2015). RH1.5D solves the non-LTE radiative transport for spectral lines in partial redistribution (PRD). This is important in the synthesis of spectral lines where there is a correlation between the frequencies of the absorbed and re-emitted radiation such that the line source function varies with frequency. RADYN assumes complete redistribution, where the frequencies of the ingoing and outgoing photons are independent. While this is reasonable for the energetics in RADYN simulations, the assumption is best suited for modelling of photospheric lines. A more accurate treatment of photon scattering is required when modelling chromospheric lines, hence PRD is generally assumed in the synthesis of Mg II h and k (Milkey & Mihalas 1974; Leenaarts et al. 2013a,b) and Ca II H and K (Vardavas & Cram 1974; Shine et al. 1975; Bjørgen et al. 2018).

The atmospheric evolution from the RADYN simulations was used as input to the RH code at every second. We added a microturbulence term of 7 km s^{-1} to produce a plage-like chromosphere following the work of Carlsson et al. (2015). This atmospheric structure is not proposed as a realistic model of plage, but aims to reproduce the observed average Mg II k line profile. However, we note that the atmosphere still fails to reproduce the single peaked profiles observed in plage.

In the synthesis of hydrogen and calcium, we used the 5 level-plus-continuum H I and Ca II atoms that are standard in RH1.5D, while for magnesium, we used the 10 level-plus-continuum Mg II atom from Leenaarts et al. (2013a). All lines were treated in PRD, even though Bjørgen et al. (2018) found that it is less important for Ca II 854.2 nm. The synthetic spectra were calculated for the first 40 s, allowing us to investigate the chromospheric emission during the heating and post-heating phases in detail.

Testa et al. (2020) investigated the Mg II 279.882 nm triplet in a variety of RADYN simulations, including the empty loop models presented here. The Mg II triplet lines (279.160 nm, 279.875 nm and 279.882 nm) are located in the wings of the h and k lines, and are expected to be in emission when large temperature increases are present in the lower chromosphere (Pereira et al. 2015). We included the Mg II 279.882 nm triplet in the formation height analysis, in order to compare it to the Ca II 854.2 nm line which is formed in the lower chromosphere.

2.4. Contribution function to the line intensity

The final part of the spectral diagnostics consisted of analysing the contribution function to the emergent intensity. Following Carlsson & Stein (1997), the contribution function was calculated as

$$C_{I_\nu}(z) \equiv \frac{dI_\nu(z)}{dz} = S_\nu \tau_\nu e^{-\tau_\nu} \frac{\chi_\nu}{\tau_\nu}. \quad (2)$$

The first term on the right-hand side gives the total source function S_ν , which is dependent on frequency because we assume PRD. The optical depth factor $\tau_\nu e^{-\tau_\nu}$ represents the Eddington-Barbier part of the contribution function, and peaks at $\tau_\nu = 1$. The χ_ν/τ_ν term is the ratio of the opacity over optical depth which is sensitive to velocity gradients in the atmosphere, and is responsible for the line asymmetries. The factor is dominant when the opacity is large at small optical depths, which typically occurs when the velocity gradients are strong.

The opacity χ_ν and source function S_ν were calculated by the RH1.5D code for selected wavelengths, and the optical depth τ_ν was calculated by integrating the opacity over the entire loop length. The contribution function to the line intensity was calculated in the range $z = 0\text{--}2$ Mm, which is the part of the loop that has near-vertical geometry.

3. Results

3.1. Atmospheric response to non-thermal electrons

Figures 2 and 3 show the evolution of temperature T , vertical velocity v_z , electron number density n_e , and beam heating rate H for the loops with $T_{LT} = 1$ and 3 MK. The quantities are plotted in the range of interest in terms of formation heights for the relevant spectra. For the temperature, velocity, and density, we focus on the first 20 s of the simulations in order to investigate the heating and early relaxation phase in more detail. The beam heating rate is shown for the first 10 s, corresponding to the loop heating time. These figures essentially reproduce figures 2 and 3 of Polito et al. (2018), but show the evolution in the lower 3 Mm of the loops in more detail. We refer to Polito et al. (2018) for more details on the evolution of the entire 15 Mm loop-length.

3.1.1. Empty loop

The difference between heating models is largest in the loop with initial low electron number density and 1 MK loop-top temperature (see Fig. 2). During the 10 s heating phase, electrons with $E_C = 5$ keV deposit most of their energy in the corona and TR, leading to an increase in temperature and density of the ambient plasma of about one order of magnitude. This causes the TR to shift to lower depths, as can be seen in panel (a). The large (negative) velocities in panel (d) result from an increase in pressure due to the rise in temperature. The upward motion of hot plasma ('chromospheric evaporation') increases with height, and as a consequence, the loops are filled with hot plasma and the electrons start depositing energy at increasing heights (see panel (j)). The positive velocities are indicative of a significant downflow of plasma towards the chromosphere, caused by the condensation of hot plasma filling the flaring loop.

The heating models with $E_C = 10$ and 15 keV (middle and right columns) have more energetic electrons in their distributions. Higher-energy electrons can travel through the corona depositing only a small fraction of their energy, allowing them to penetrate deeper into the atmosphere. Panels (h) and (i) show that the deposited energy is at a maximum between $z \approx 1.5$ and 2.1 Mm within the first few seconds. The radiative losses from ionised hydrogen and helium are not able to balance the large amount of energy deposited, resulting in a temperature increase in this region (as can be seen in panels (b) and (c)). The up-flowing velocities in panels (e) and (f) are lower compared to the 5 keV model. This is due to the difference in sites where electron energy is deposited. While the 5 keV electrons deposit most of their energy in the corona and TR, the 10 and 15 keV

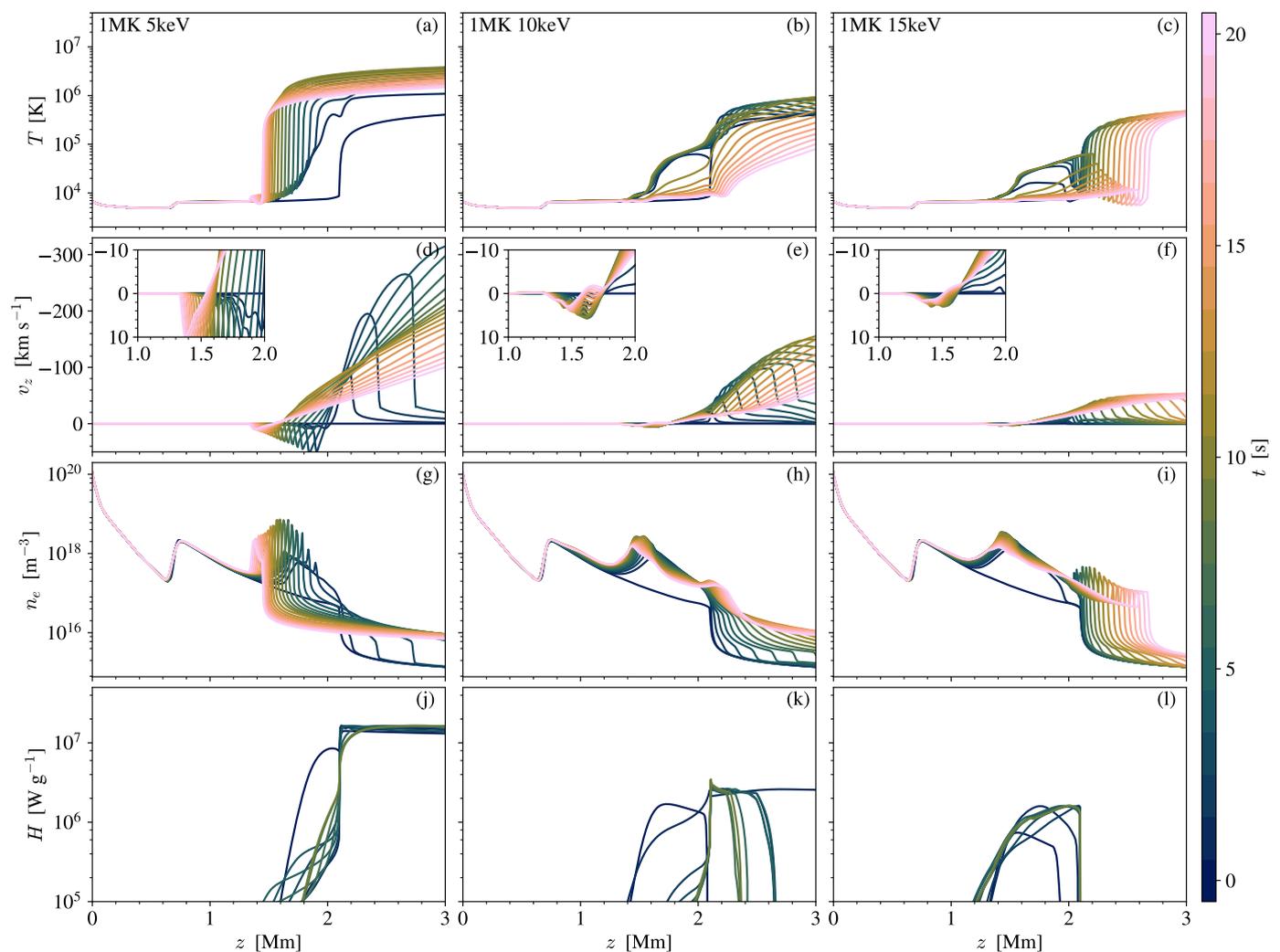


Fig. 2. Atmospheric response to the RADYN simulation for the 15 Mm half-loop with initial apex temperature of 1 MK. Each column corresponds to the heating model with $E_C = 5, 10,$ and 15 keV, and each row represents the temperature, velocity, electron number density, and beam heating rate. Negative (positive) velocities correspond to upflows (downflows). The quantities are plotted in the range $z \in [0, 3]$ Mm at 1 s intervals. Temperature, velocity, and electron number density are displayed for the first 20 s of the simulations, while the beam heating rate is given for the duration of the injected electron beam (10 s). The insets in panels (d)–(f) show the velocity in the region $z = 1$ – 2 Mm saturated to ± 10 km s $^{-1}$.

electrons mostly deposit their energy in the chromosphere where the plasma is denser. Dense plasma is harder to accelerate, and the combination of this and the fact that energy radiates away more efficiently contributes to weaker upflowing velocities. Consequently, the upflows cause the TR to move towards greater heights.

The upward motion of hot plasma in the 10 and 15 keV heating models causes the density of the loops to increase over time, and similarly to the 5 keV model, the electrons start to deposit energy at increasing heights (see panels (k) and (l)). However, we note that in the 15 keV model the increase in height where energy is deposited does not change much over the 10 s heating duration compared to the other models. The 5 keV electrons are more effective at heating the corona, and the large upflowing velocity demonstrates that this heating model is more efficient at filling the flaring loop with high-temperature plasma.

3.1.2. Pre-heated loop

Figure 3 shows the atmospheric response to the heating models for the loops with initial apex temperature of 3 MK. As men-

tioned earlier, the hotter and denser loops represent strands that have been previously heated. Therefore, the initial atmospheres have transition regions that are located at lower depths compared to the 1 MK loop models. Panels (a)–(c) show a temperature increase in the corona for all E_C values. Due to the higher initial density of the loops, the injected electrons deposit more of their energy in the corona. In the 5 keV heating model, the energy deposited directly into the corona (as indicated by the increasing heating rate with height in panel (j), but also shown in Polito et al. (2018)) causes the temperature and pressure to increase, leading to both upflows and downflows of plasma. However, not all electrons deposit their energy in the corona. Some of the electrons are able to reach the TR, where they deposit their energy at $z \approx 1.5$ Mm. The local heating of the TR in combination with the conduction front from the corona results in a pressure gradient that produces both upward motion of plasma towards the corona as well as downflowing plasma into the TR.

The intermediate- and high-energy electrons initially deposit their energy in the TR, resulting in upflows of hot plasma into the corona. As the loops are filled with high-temperature plasma, the electrons start to deposit energy at greater heights, as can be seen

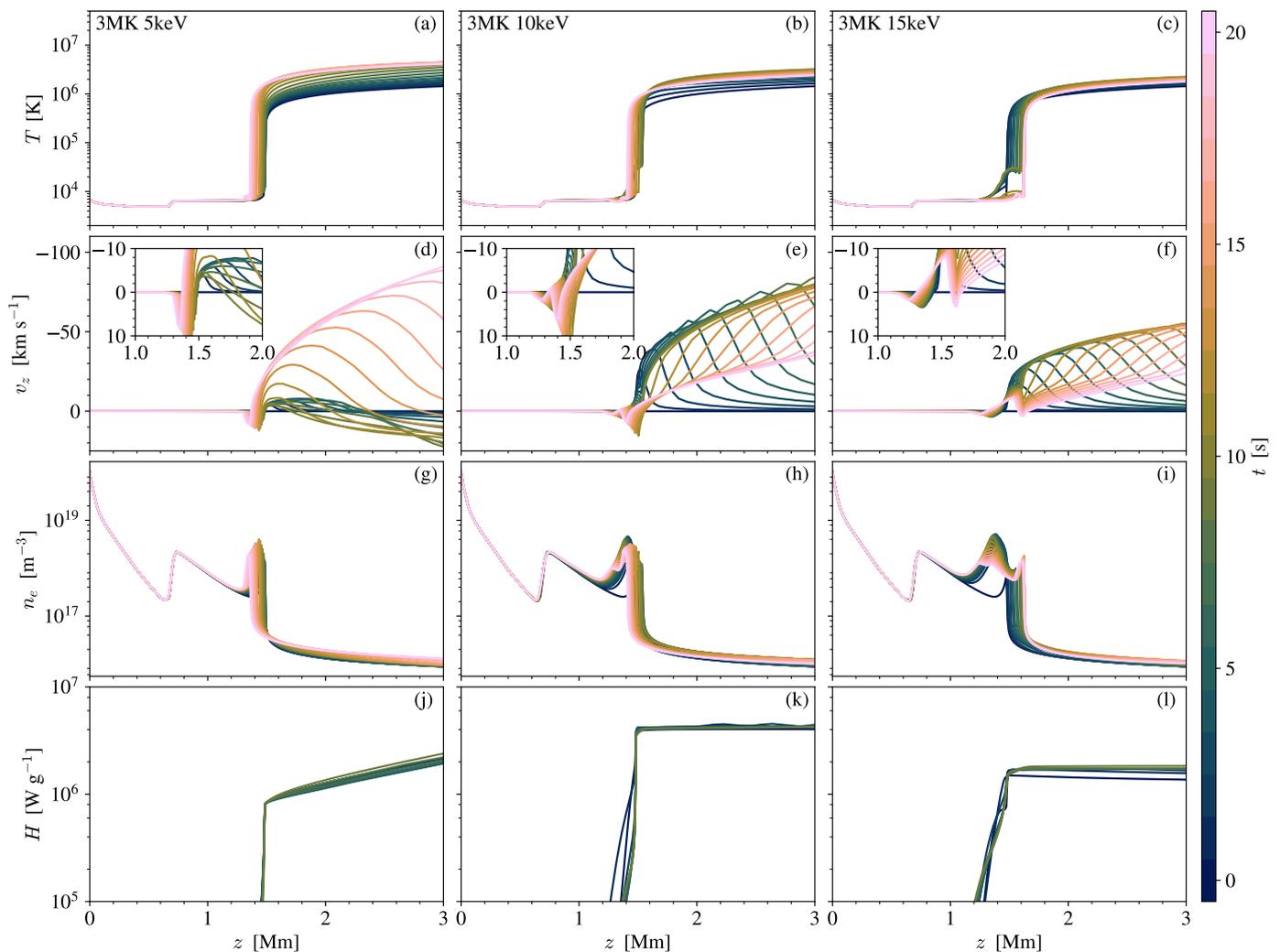


Fig. 3. Atmospheric response to the RADYN simulation for the 15 Mm half-loop with initial apex temperature of 3 MK. See Fig. 2 for an explanation of the different panels.

in panels (k) and (l). Panels (e) and (f) show that as the height where energy is deposited increases, the evaporation front travels upwards along the loop. The moving TR region in panel (c) shows resemblance to the 10 and 15 keV runs in Fig. 2, where the properties of dense plasma in combination with direct heating results in an increased temperature. However, the atmospheric response is much slower for the 3 MK loop because the plasma is harder to accelerate, hence the TR does not move much during the first 20 s.

3.2. Chromospheric emission

Figures 4 and 5 represent the time evolution of the synthetic Mg II h, Ca II K, and Ca II 854.2 nm spectra, along with their light curves, from the different heating models. The light curves are obtained by integrating the emission over the spectra in the wavelength range for each timestep. We note that the Mg II h results are similar to Mg II k, and that the Ca II K results are similar to Ca II H. Therefore, we only present our findings from the Mg II h and Ca II K lines. Polito et al. (2018) analysed the Mg II h line in detail, and we reproduce the Mg II h results in order to make a direct comparison with the Ca II lines and study the effect from energy deposited by non-thermal electrons.

Figure 4 shows the chromospheric emission for the 1 MK loop, where each column represents the different cutoff energies E_C . The three spectral lines in the $E_C = 5$ keV model are behaving similarly, but with some variation. The Mg II h and Ca II K profiles have redshifted components that are caused by the downflow of cool plasma into the chromosphere, as seen in Fig. 2 (d). The Mg II h line also has a blueshifted component, and is characterised by multiple peaks that are Doppler shifted up to approximately ± 50 km s⁻¹ during the heating phase. The Ca II K line shows multiple peaks in the redshifted component that are Doppler shifted to around $+40$ km s⁻¹. The Ca II 854.2 nm line shows weak emission of the red wing at the end of the heating phase. The blue peak gets broader and has a higher intensity level during the post-heating phase, as a response to the upflow of plasma into the corona and TR.

During the first 10 s of the simulations, while the electrons are still injected, the profiles in the $E_C = 10$ and 15 keV models are mostly double peaked. The spectra experience an increase in line core intensity during the heating phase. This is because the high-energy electrons are able to penetrate deeper into the atmosphere before depositing their energy, as can be seen from panels (h) and (i) in Fig. 2. During 5–10 s, the Mg II h line in the 15 keV model shows faint emission of the blue wing, which is Doppler shifted to approximately -30 km s⁻¹. After the heating

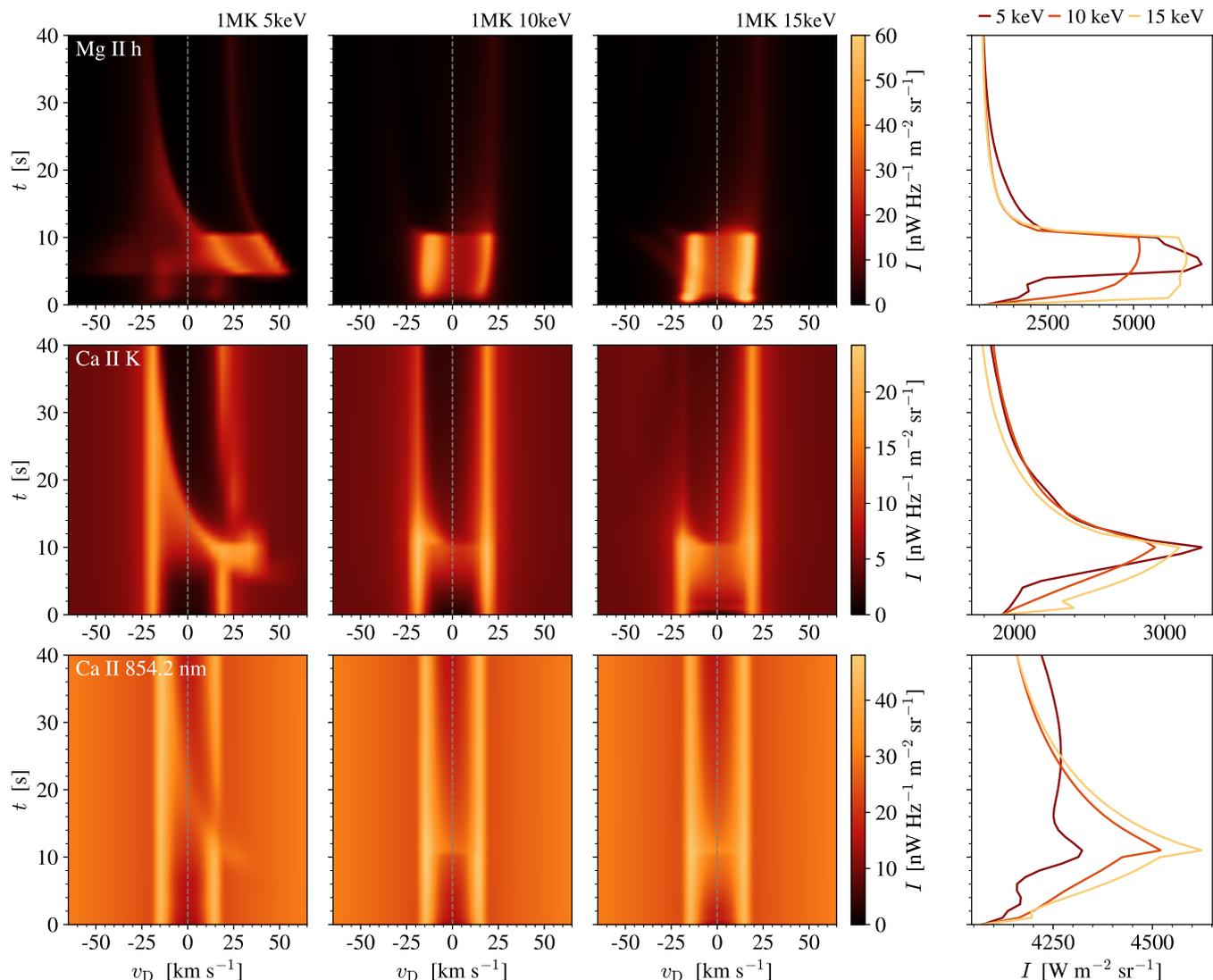


Fig. 4. Spectral evolution of Mg II h, Ca II K, and Ca II 854.2 nm spectra for the heating models with $T_{LT} = 1$ MK and $E_C = 5, 10,$ and 15 keV. The x -axes in the first three panels of each row are in units of Doppler offset, where negative (positive) velocities indicate blueshifts (redshifts). The panels of the right column show light curves from each heating model, where the x -axes are in units of integrated intensity. The first 40 s of the simulations are shown, including the 10 s electron injection phase. The Mg II h spectra are clipped at $60 \text{ nW Hz}^{-1} \text{ m}^{-2} \text{ sr}^{-1}$ to emphasise the less bright features. Line profiles for $t = 0, 7,$ and 15 s are shown in Fig. 6.

phase, Mg II h and Ca II K are dominated by a red peak, and the profiles become almost single peaked. This is apparent from the solid pink lines in Fig. 6.

The light curves from the Mg II h spectra show a sudden increase in intensity as a response to the electron beams. The profile from the 5 keV model is generally less intense, but the integrated emission increases as a result of the broadening of the line. The light curves from the Ca II spectra have a slower increase and peak around 10 s. While most of the light curves decrease after the heating is turned off, the Ca II 854.2 nm integrated emission starts to increase around 15 s. This is due to the broadening and increased emission of the blue peak.

Figure 5 shows the spectra for the denser 3 MK loop. The 5 keV simulation shows qualitatively similar effects on the Mg II h and Ca II K lines, while the effect on the Ca II 854.2 nm line is less significant. The Mg II h and Ca II K lines both show broadening of the blue peak at around 10 s. During the post-heating phase, the entire Mg II h line is shifted to the red before

slowly reverting back. The blue component of Ca II K is characterised by double peaks, and the general behaviour of the intensity after the heating is turned off is similar to that of Ca II K in the 5 keV empty loop model. Due to the density of the loop, the electron energy is dissipated directly in the corona and TR. The chromosphere is not heated significantly, and the effects seen in Mg II h and Ca II K are due to local plasma flows.

In the 10 and 15 keV simulations, there is an increase in intensity as a function of E_C that is particularly apparent in Mg II. The peak intensity of Mg II h in the 10 keV model is comparable to that of an empty loop. However, the energy deposited from the 15 keV electrons affect the chromospheric line greatly, where the intensity of the red peak is much stronger compared to the empty loop (see Fig. 6). The Ca II 854.2 nm profiles in the $E_C \geq 10$ keV simulations are comparable to the results of the empty loop, both in line features and evolution of the light curves. The Ca II K profiles are greatly affected by the intermediate- and high-energy electrons, and are characterised by multiple peaks and strong

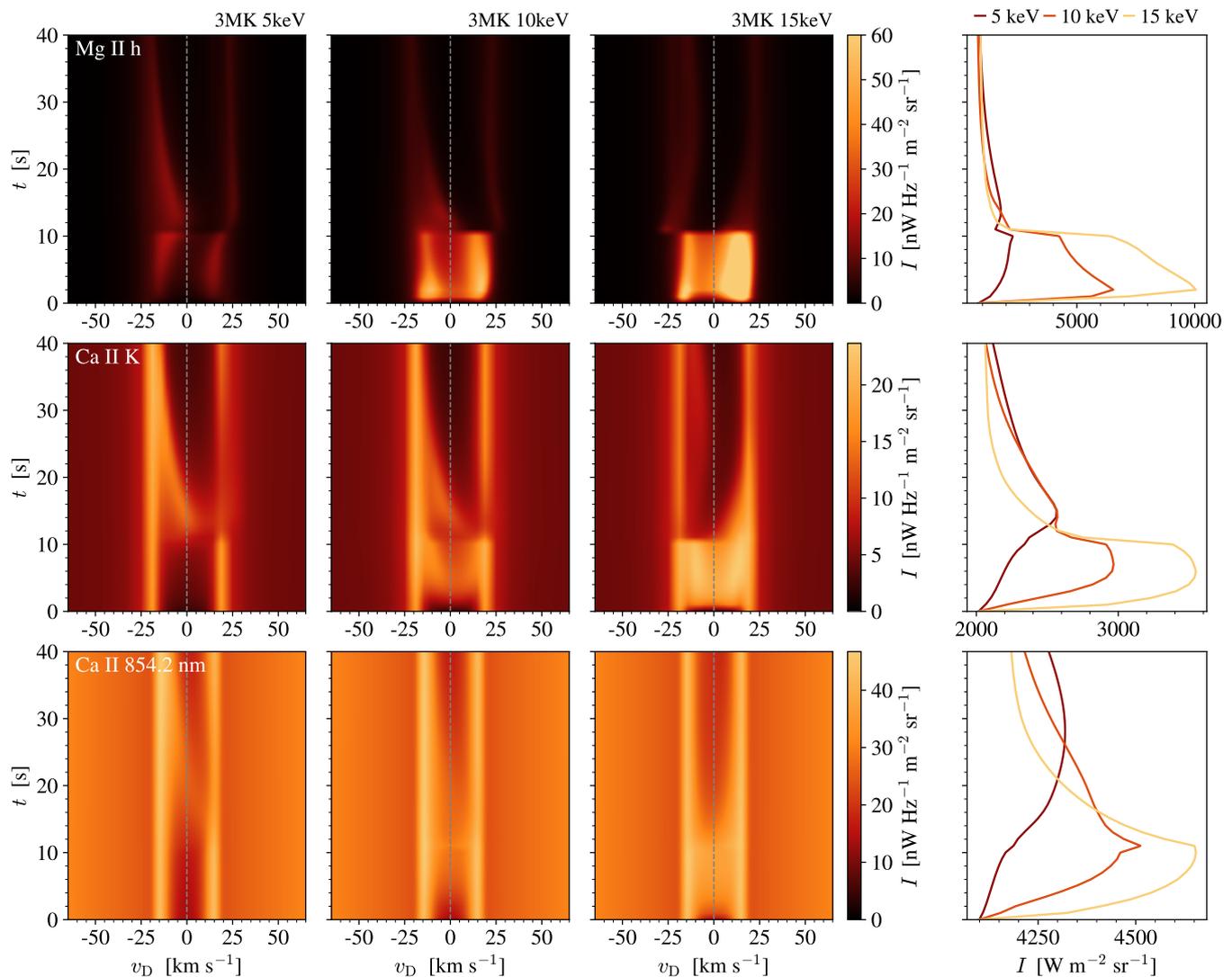


Fig. 5. Synthetic spectra for the $T_{LT} = 3$ MK heating models. See the caption of Fig. 4 for more details.

emission of the line core. The latter is reflected in the light curves, which are similar to the Mg II integrated emission from the 10 and 15 keV models in that they have a steep decline.

Figure 6 shows the Mg II and Ca II line profiles to better visualise the shapes and features that arise during the simulations. Particularly, the faint emission of the blue wing of Mg II h in the 15 keV empty loop model is easier to observe (see the solid green line in the upper right panel). In the 5 keV empty loop model and the 5 and 10 keV dense loop models, the Ca II K core is shifted to the red after the heating is turned off, and the blue component is double peaked. The Ca II 854.2 nm profiles show that the response to the non-thermal electrons is mostly changing the absorption feature, which gets brighter as time progresses. In the 5 keV models, the red emission peak of Ca II 854.2 nm is weaker at later timesteps.

Figure 7 represents the synthetic H α and H β lines from all the simulations, and are given for the same timesteps as in Fig. 6. The initial profiles are almost identical across the simulation span, where both H α and H β are in absorption. In the 1 MK low-energy model, both spectra are in emission during the heating phase. The H α line is characterised by a redshifted component with multiple peaks, while the redshifted H β peak is severely

broadened. In the post-heating phase, both lines are back in absorption, where their cores are Doppler shifted to approximately +15 km s $^{-1}$. In the 3 MK loop, both lines are in absorption during the heating and post-heating phases, showing no significant effects as a consequence to the nanoflare event.

The intermediate- and high-energy models (both 1 and 3 MK loops) show similar effects in the H α and H β lines. During the heating phase, the spectra are double peaked. In the 10 keV models, the blue peak is dominating the spectra in the empty loop, while the red peak dominates the spectra in the dense loop. In the 15 keV dense loop model, the red peak of both lines is more pronounced. The latter is also true for the H α line in the 1 MK loop, while H β is almost symmetric around the line core. All heating models show that the line cores are in absorption during the post-heating phase, but some of the line cores are either redshifted or blueshifted as a response to the non-thermal electrons.

3.3. Formation of line intensity

Figures 8 and 9 show four 2×2 diagrams of the intensity formation of Mg II h and Ca II K from two of the heating models presented in this work. Because of the extensive parameter sur-

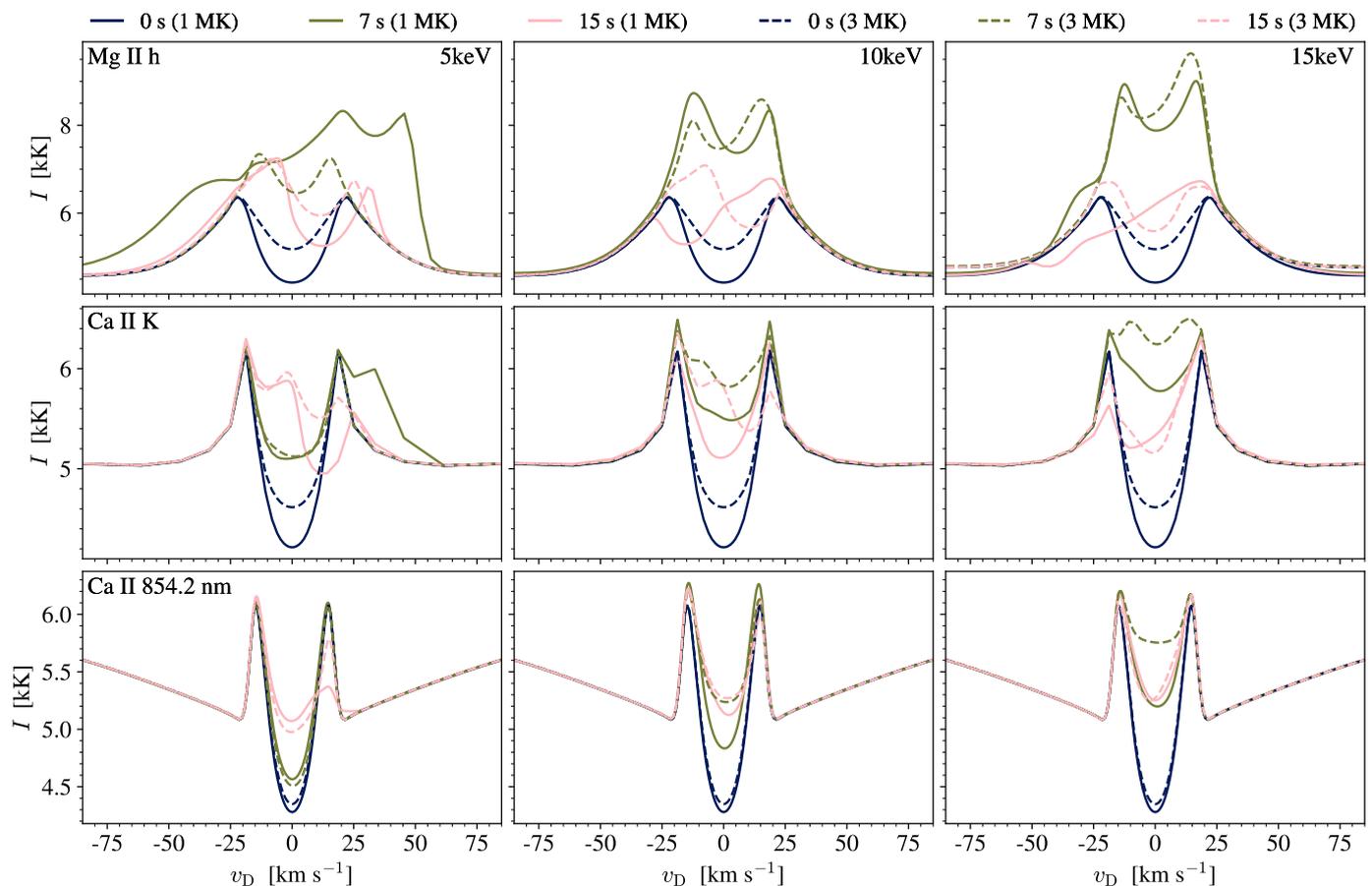


Fig. 6. Synthetic Mg II h, Ca II K, and Ca II 854.2 nm line profiles. The solid and dotted lines represent the 1 and 3 MK loops, respectively. Line profiles during the pre-heating ($t = 0$ s), heating ($t = 7$ s), and post-heating ($t = 15$ s) phases are shown.

vey, we selected two distinctively different models. We present the Mg II h and Ca II K spectra subject to low-energy (5 keV) electrons in the 1 MK loop and high-energy (15 keV) electrons in the 3 MK loop. The specifics of the simulations are given in the top left corner of each diagram. The panels in each subfigure represent the terms in Eq. 2, and are strongly inspired by the four-panel diagrams in Carlsson & Stein (1997).

Figure 8 (a) shows the Mg II h line during the heating phase in the 5 keV empty loop model. The downward velocity gradient is reflected in the χ_v/τ_v term, exhibiting a strong variation in wavelength. As a result, the $\tau_v = 1$ height of the red peak is shifted to larger Doppler offsets. The source function is equal to the Planck function from the photosphere up to around 1.2 Mm where it decouples. Since the Planck function is given in units of brightness temperature, it also represents the temperature itself. The strong rise in temperature at 1.7 Mm gives rise to a significant increase in the source function. The peak in S_v is further responsible for the multiple peaks observed in the redshifted component of Mg II h. Investigation of the $\tau_v = 1$ height reveals that what appeared to be a double peaked blue component is emission of the blue wing. This is further supported by the increase in S_v around the blue wing. The contribution function is strongly dominated by the downward velocity gradient because of its dependence on χ_v/τ_v . Signatures of the upflows causing emission in the blue wing are therefore difficult to enhance visually. We have solved this by applying gamma correction to χ_v/τ_v and C_I , simply allowing us to highlight the darker regions.

The intensity formation of Mg II h during the post-heating phase is given in Fig. 8 (b). The velocity gradient from the down-flowing plasma is weaker compared to the heating phase, and has less impact on the χ_v/τ_v term. The source function is coupled to the Planck function until approximately $z = 1.4$ Mm. The sudden increase in temperature at 1.5 Mm causes the source function to peak, giving rise to two emission peaks in the intensity profile. The increase in temperature and source function gives rise to an additional peak at approximately 0 km s^{-1} , causing the minimum of the central absorption to shift to the red.

Figure 8 (c) shows the intensity formation of Ca II K during the heating phase. The Ca II K line is symmetric between $\pm 25 \text{ km s}^{-1}$, as the $\tau_v = 1$ curve is not affected by the strong velocity gradient in this range. Here, the shape of the profile is mostly defined by the source function, which is coupled to the Planck function up to 0.9 Mm. At around $+35 \text{ km s}^{-1}$, the $\tau_v = 1$ height peaks as a response to the sudden increase in S_v , which is caused by the increase in temperature due to the strong down-flowing velocity. This gives rise to a separate redshifted component at $+35 \text{ km s}^{-1}$ in the Ca II K line. We note that PRD affects Ca II K less than Mg II h which becomes apparent in the source function panel in the top right: there is little variation in S_v along the x -axis for most of the atmosphere and S_v appears like a horizontal grey band.

Panel (d) in Fig. 8 gives the formation of Ca II K during the post-heating phase. Similar to Mg II h (see panel (b)), the less strong downflowing velocity is reflected in the χ_v/τ_v term, showing the largest increase around 0 km s^{-1} . The sudden increase in

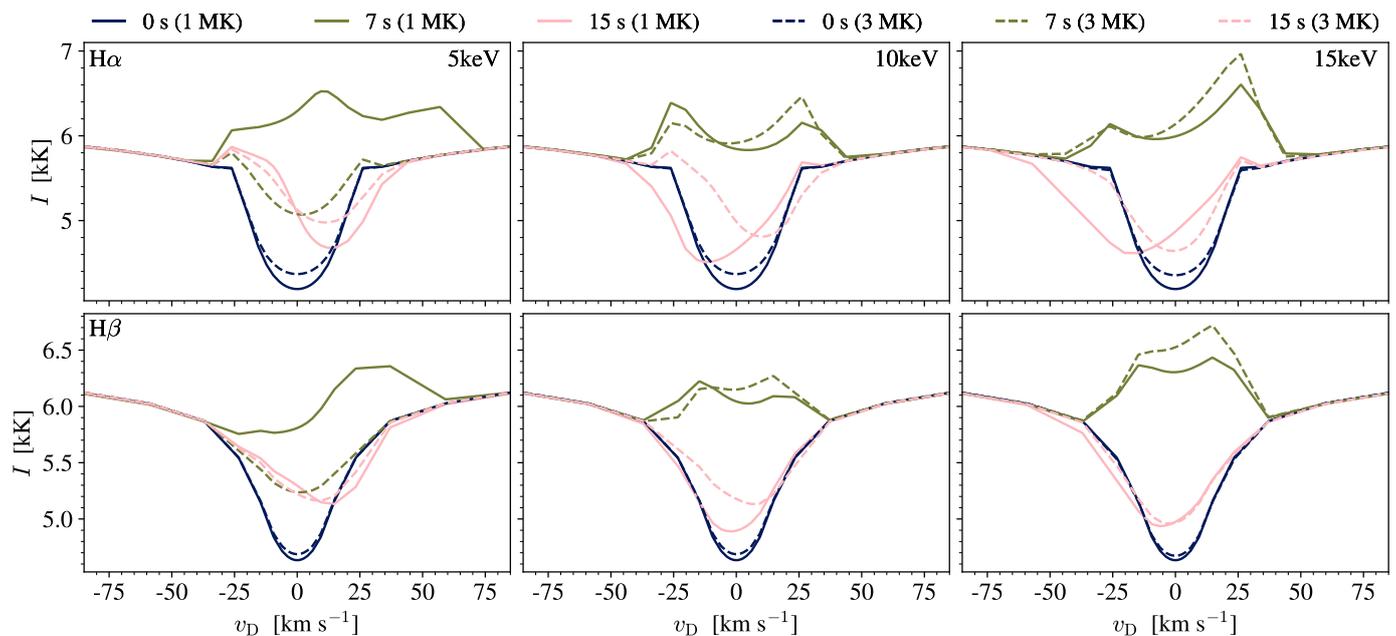


Fig. 7. Synthetic H α and H β spectra. See the caption of Fig. 6 for more details on the different panels.

temperature at 1.5 Mm causes the source function to peak, giving rise to the peak near the centre of the line at approximately $+7 \text{ km s}^{-1}$. The intensity of the peak at -20 km s^{-1} is stronger than the peak at $+35 \text{ km s}^{-1}$ as a consequence of the formation of the peak near the line centre. We also note that the contribution function peaks at different Doppler offsets for Mg II h (0 km s^{-1}) and Ca II K ($+7 \text{ km s}^{-1}$), even though the lines form at the same height. This shows that the velocity is not the only contributing factor to the formation of spectral features, and this underlines the risk of simply interpreting the Doppler offset of spectral features as a direct measure of plasma flows in the atmosphere.

Figure 9 (a) shows the formation of Mg II h during the heating phase in the 15 keV dense loop model. The source function is coupled to the Planck function up until 1.2 Mm, but continues to follow it closely up to around 1.4 Mm. The peak in S_ν gives a larger contribution to C_{I_ν} around the red component. As a consequence, the red peak experiences a larger intensity increase and the minimum of the central absorption is shifted to the blue.

The intensity formation of Mg II h during the post-heating phase is shown in Fig. 9 (b). The asymmetric $\tau_\nu = 1$ curve is caused by the variation in velocity between $z = 1.3\text{--}1.6 \text{ Mm}$. The source function is coupled to the Planck function up to 1.3 Mm. From this height, the source function follows a flat slope before it suddenly decreases at 1.6 Mm. The contribution function shows that the peaks of Mg II h form at different heights. The line profile is slightly shifted to the blue, with the minimum of the central absorption at approximately -5 km s^{-1} . The blueshift is caused by the general upflow pushing the TR towards greater heights (see Fig. 3 (c)).

Figure 9 (c) shows the intensity formation of Ca II K during the heating phase. The double peaks in the blue component are caused by the velocity gradient of the upflowing plasma. The source function peaks as a result of the temperature increase, but the upward motion of plasma dilutes the signal such that S_ν is weaker around the $\tau_\nu = 1$ height of the line core. Still, the contribution to the line core is large enough to increase the intensity, making the emission peaks less pronounced.

Figure 9 (d) shows the Ca II K line during the post-heating phase. The varying velocity of the fluid causes the asymmetric shape of the optical depth unity curve, where the $\tau_\nu = 1$ height increases up to 1.6 Mm at negative Doppler offsets. Similar to the Mg II h line, Ca II K is blueshifted to around -5 km s^{-1} . The source function and Planck function are coupled to approximately 0.9 Mm, but does not differ significantly until around 1.4 Mm. As the Planck function (and temperature) increases, the source function decreases. As a result, the source function makes a stronger contribution to C_{I_ν} at lower heights. This is reflected in the intensity profile, where the red peak is more intense than the blue peak.

Figure 10 shows the formation heights of the Mg II h, Mg II 279.88 nm triplet, Ca II K, Ca II 854.2 nm, H α , and H β line cores. We include the Mg II triplet to investigate how the formation height of another deep chromospheric line compares to Ca II 854.2 nm. Since the lines are formed under optically thick conditions, we define the formation height to be the maximum height of the $\tau_\nu = 1$ surface. During the heating phase in the 1 MK heating models (top panels), the Mg II h and Ca II cores are formed at lower heights compared to their initial $\tau_\nu = 1$ height, while the Mg II triplet, H α , and H β cores are formed higher in the chromosphere compared to their initial $\tau_\nu = 1$ height. In the 5 keV model, there is a sudden increase in the height of formation for all line cores. This sudden increase is progressively later for spectral lines with lower formation heights, which is caused by local effects from the receding TR (see Fig. 2 (a)). As a response to the moving TR, the line cores form over a narrow region in the cooling phase due to the compression of the chromosphere. The formation heights of the line cores in the 10 and 15 keV models increase at 10 s as a result of the sudden decrease in temperature when the heating is turned off. As the temperature continues to decrease, the $\tau_\nu = 1$ heights approach their pre-heating values.

The bottom panels of Fig. 10 represent the 3 MK dense loop models. Since the loop is pre-heated, the TR is located at lower heights compared to the empty loop. The line cores therefore form at lower heights as well. The 5 keV electrons do not af-

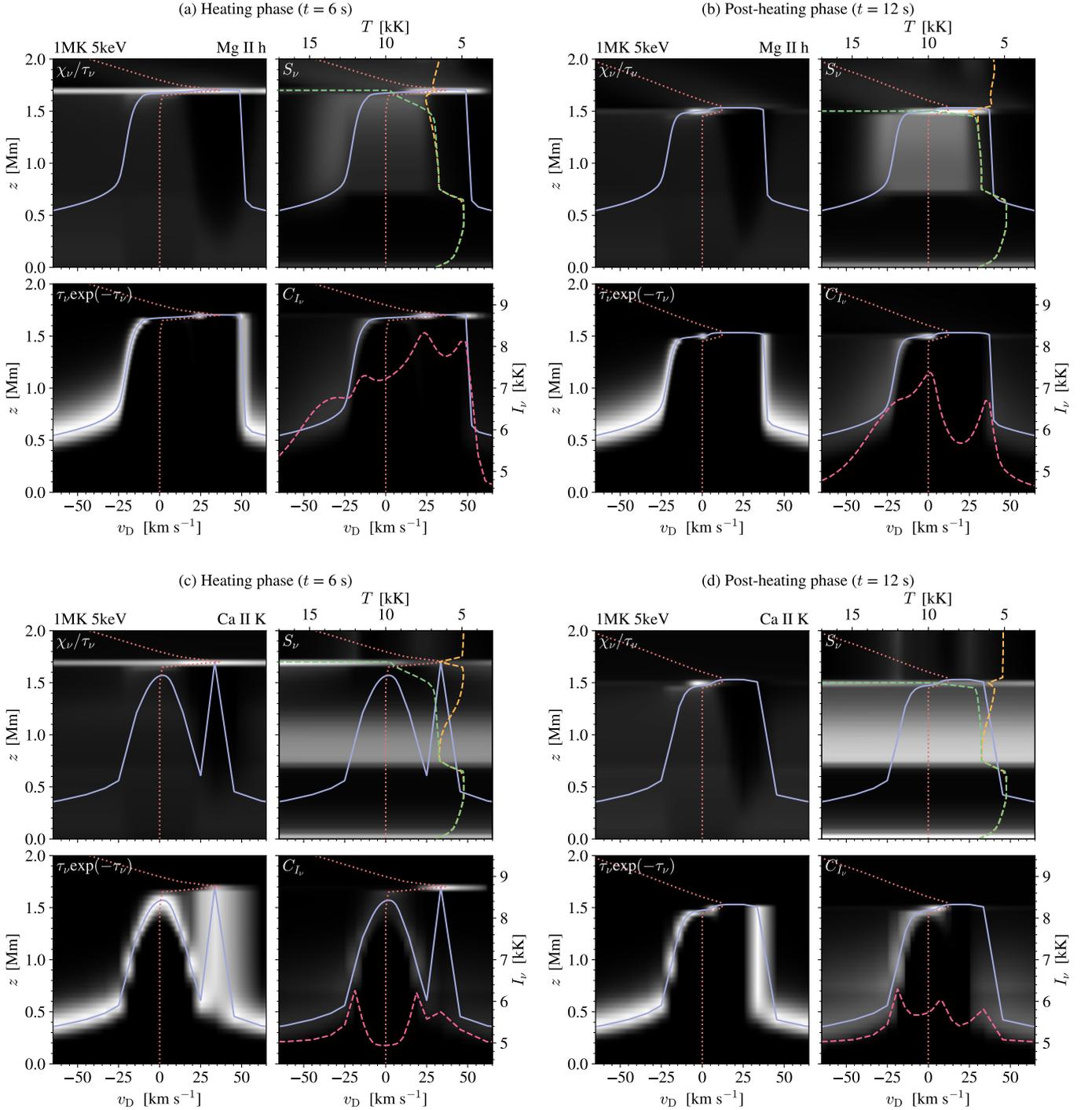


Fig. 8. Intensity formation of the Mg II h (top) and Ca II K (bottom) spectral lines from the 5 keV empty loop model. The subfigures show the heating phase ((a) and (c)) and early post-heating phase ((b) and (d)) for the two lines. The quantities given at the top left corner are shown in greyscale as functions of frequency from line centre (in units of Doppler offset) and height z . The $\tau_v = 1$ height (purple) and vertical velocity (red dotted) are displayed in all panels. Negative (positive) velocities correspond to upflows (downflows). In addition, the top right panels show the total source function at $v_D = 0$ (yellow dashed) and Planck function (green dashed) in units of brightness temperature, with high values to the left. The bottom right panels also contain the intensity profile (pink) as brightness temperature. Gamma correction is added to the χ_v/τ_v and C_{I_v} terms to amplify the weaker signals.

fect Mg II h, Ca II K, Ca II 854.2 nm, and H α significantly, hence these line cores form at similar heights during the heating phase. The chromosphere is affected by the intermediate- and high-energy electrons, but the response is slow since the loop is already heated. As a result, the formation heights of the Mg II h

and Ca II K cores do not drop as significantly within the first few seconds as in the empty loop. In the 10 keV model, the formation heights decrease as the TR recedes to lower depths. In the 15 keV model, the increase in $z(\tau_v = 1)$ after the heating is turned off is a result of the TR moving to greater heights.

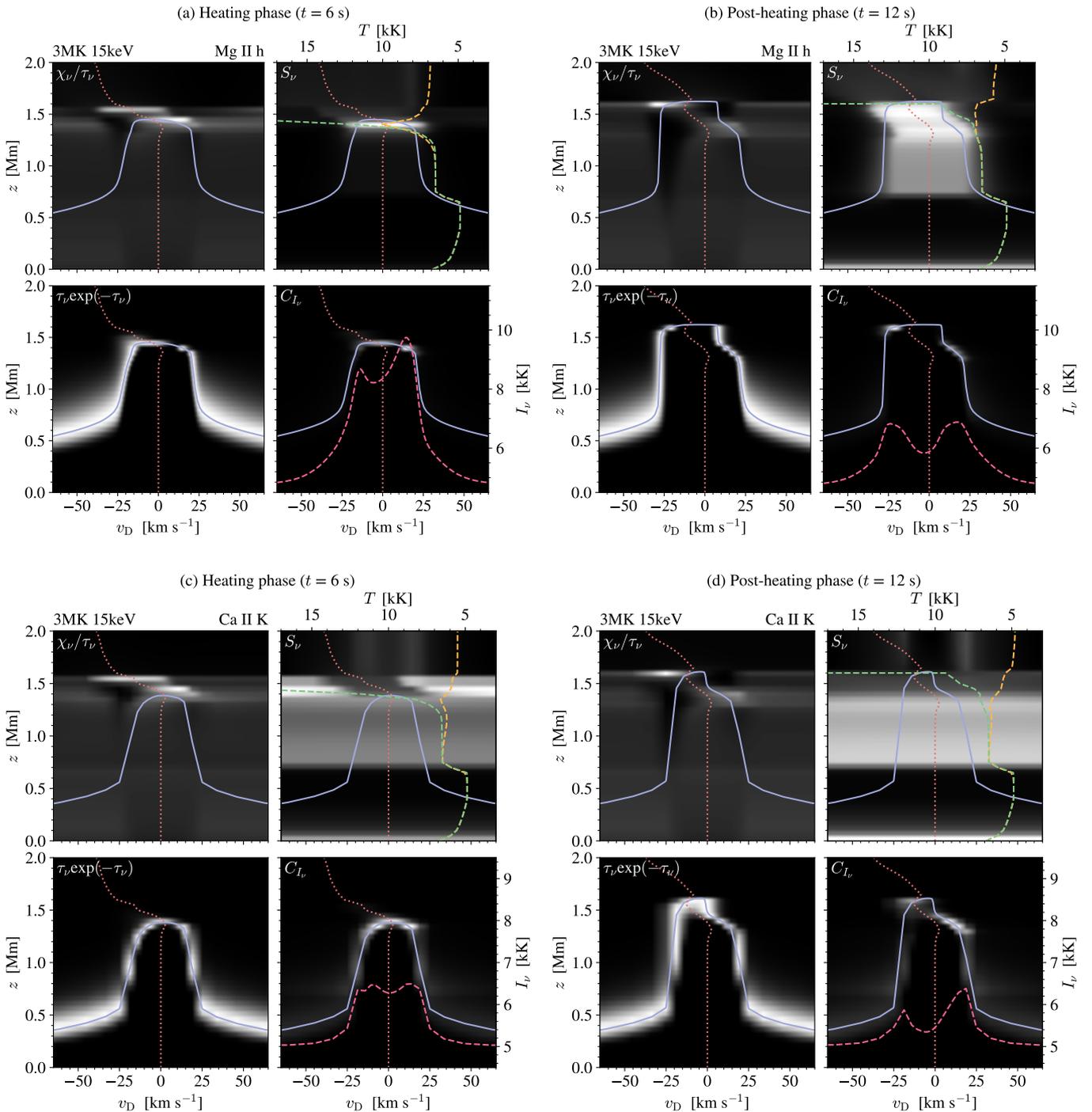


Fig. 9. Intensity formation of the Mg II h (top) and Ca II K (bottom) spectral lines from the 15 keV dense loop model. We note the different scaling of the brightness temperature in the top and bottom panels. See the caption of Fig. 8 for more details.

The increase in formation height during the heating phase for the Mg II triplet and H β cores shows that these lines are sensitive to heating. This result was also found by Testa et al. (2020) for the Mg II triplet. The 5 keV dense loop model even demonstrates that minor atmospheric responses from energy deposited by the non-thermal electrons affect their $\tau_\nu = 1$ heights. In addition, the panels show that Ca II 854.2 nm is the deepest diagnostic when electrons are being injected into the loop.

4. Discussion

By extending the spectral line diagnostics of Polito et al. (2018) and Testa et al. (2020), we find that also dominant chromospheric lines in the visible are affected by coronal nanoflare events. The signatures that arise from the non-thermal electron beams highly depend on the initial physical conditions of the loop in combination with the low-energy cutoff E_C of the power-law distribution of accelerated electrons. Most importantly, the distance the electrons travel before their energy is deposited is largely dependent on the density of the loop and the energy of the electrons. In

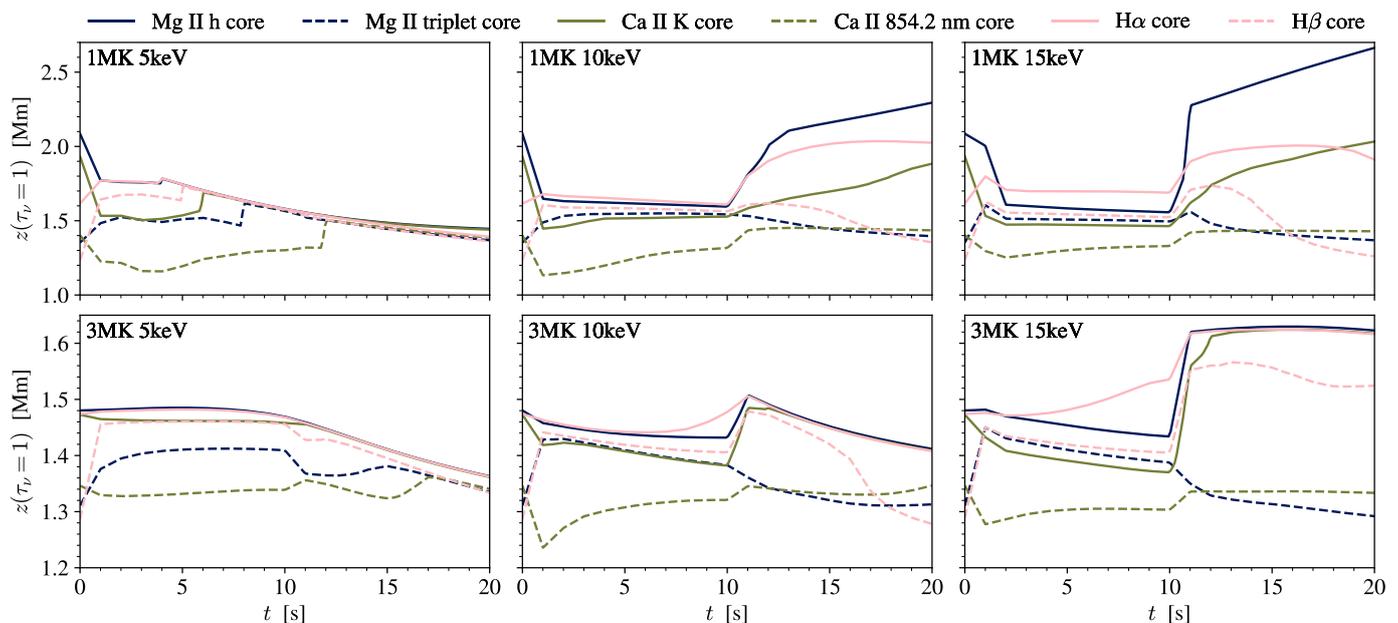


Fig. 10. Maximum formation height of the Mg II h, Mg II 279.882 nm triplet, Ca II K, Ca II 854.2 nm, H α , and H β line cores. The first 20 s of the simulations are shown, where time is given on the x -axis. Each panel represents the different heating models, as stated in the upper left corners.

low-density loops, there are fewer collisions with the ambient plasma, and vice versa. The spectral line signatures are not only affected by the energy deposited directly in the regions where they form, but also by the response to the effects from electron beam energy dissipated elsewhere in the atmosphere, for example plasma flows.

The most significant effects are seen in the empty loop, where the atmospheric response to $E_C = 5$ keV electrons causes complex line profiles that are redshifted to large Doppler velocities. Through investigation of the intensity formation, we have shown that the spectral lines are highly affected by the plasma flows through the χ_ν/τ_ν term. The increase in coronal and TR temperatures leads to the strong downflows responsible for the redshifts, and it is a result of the electrons depositing most of their energy in those regions. The large upflows (negative velocities in Fig. 2 (d)) are in agreement with the findings by Testa et al. (2014) and Reep et al. (2015), who showed that low-energy electrons could drive ‘explosive evaporation’. The atmospheric response to 5 keV electrons strongly depends on the sites where the energy is deposited.

The stopping depth of the electrons in the $E_C = 10$ and 15 keV energy distributions is deeper into the atmosphere, mainly in the TR and chromosphere. However, we note that the 10 keV electrons in the 3 MK loop deposit more energy in the corona and TR, similarly to the 5 keV models. A portion of the electrons is able to travel to the chromosphere, where the increase in chromospheric emission during the heating phase is significantly larger than in the 5 keV models. The height where the energy is deposited strongly depends on the density of the loop, which becomes apparent when comparing the 10 keV empty and dense loop models. The increased density causes significantly more electrons from the distribution to collide with the ambient plasma at great heights. The atmospheric response to the 10 keV electrons in the dense loop therefore resembles that of the 5 keV models, while the response to intermediate-energy electrons in the empty loop is more similar to the high-energy models.

A range of energies is represented in the electron energy distributions (see Fig. 1), where the input parameters are free and chosen based on observations. The advantage of 1D flare modelling with RADYN is the ability to study specific energetic events in an isolated system. The atmospheric response to the electron beams provides information on features that we can look for in more advanced simulations, such as those produced with the radiative magnetohydrodynamic (MHD) Bifrost code (Gudiksen et al. 2011). We previously introduced a Bifrost model including accelerated particles by magnetic reconnection, where the reconnection sites were found by identifying locations where the magnetic topology is not conserved (Bakke et al. 2018). Frogner et al. (2020) further developed the modelling of the accelerated particles, and found that the magnetic topology is a significant factor when determining both the amount of accelerated electrons and their energy. In this type of Bifrost models, the magnetic reconnection, and hence particle acceleration, occurs over extended regions. An advantage with such models is the more realistic 3D modelling. However, with the complex configuration of Bifrost models, it is more difficult to control the parameter space of the simulations. RADYN follows the traditional approach of modelling flares (with the injection of electrons at the top of the loop), hence we can more easily isolate the details of the event. RADYN is also relatively quick to run, which is beneficial when surveying an extensive parameter space to investigate the effects of small-scale events on a variety of spectral lines.

The signatures in the chromospheric spectra give an indication of what to look for in observations. Complex line profiles as the ones shown in Fig. 4 using the 5 keV empty loop model have not been reported in the literature so far, but might suggest heating by non-thermal electrons if observed. The inclusion of the Ca II and H I lines to the diagnostics gives potential for observing small-scale events with ground-based telescopes, such as the Swedish 1-m Solar Telescope (SST; Scharmer et al. 2003) and the Daniel K. Inouye Solar Telescope (DKIST; Rimmele et al. 2020). With the CHROMIS instrument installed at the SST, it

is possible to sample the Ca II H and K and H β lines with high spectral resolution. Ground-based telescopes generally allow for higher spatial resolution as compared to millimetre observations and (extreme-)UV diagnostics observed from space. This could possibly provide tighter constraints on the spatial dimensions of the strand widths and temporal variability. We also note that Mg II h and Ca II K appear to be more sensitive to events at the top of the chromosphere than Ca II 854.2 nm, and coordinated observations with SST and IRIS, for instance, would be beneficial to provide more constraints on nanoflares. A comparison of observations to our numerical results, as well as Bifrost electron beam simulations, will be explored in future work. Testa et al. (2020) found that Mg II profiles from RADYN nanoflare models showed similarities to what was observed at the footpoints of AR coronal loops with IRIS. They also found that the Mg II triplet is an important diagnostic for non-thermal electrons due to its sensitivity to heating in the lower atmosphere. We will therefore include the Mg II triplet in future spectral line analysis.

In this paper, we have extended the spectral line diagnostics of Polito et al. (2018) and Testa et al. (2020) to include additional chromospheric lines. We have established that non-thermal electrons from small-scale events are able to affect line features through strong plasma flows as a response to the electrons depositing their energy in the corona and TR, but also through energy deposited directly in the chromosphere. We show that the strong chromospheric lines in the visible can also be used as diagnostics of nanoflare-size events.

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